

Title: Newton's Laws of Motion and Differential Equations

Target: On completion of this worksheet you should be able to derive equations of motion of a physical system

Force

We return to the concept of force which was already covered in the worksheet E4. Some forces are just expressed as a constant. They are independent of time and position, for example, gravity $F_g = mg$. Other forces change with time, for example, $F(t) = \sin(2\pi t)$ or with position, like the Hooke's law $F_s(x) = kx$. During motion, the position of a body changes with time. That's why we can write the Hooke's law as $F_s(x(t)) = kx(t)$ to underline time dependency of the force. In classical mechanics, we mostly focus on the forces of the form $\mathbf{F} = \mathbf{F}(t, \mathbf{x}(t), \mathbf{v}(t), \mathbf{a}(t))$, which depend on time, position of the body and the time derivatives of position (velocity and acceleration).

Newton's second law for modelling motion

In the most general form, we have $\sum \mathbf{F} = m\mathbf{a}(t)$. Depending on the nature of the force, this is used to derive the equations of motion for a physical object. For example, for a damped harmonic oscillator, the forces acting on the body are the Hooke's force $F_s = -kx(t)$, and the linear dampening term $F_d = -cv(t)$, where we wrote $x = x(t), v = v(t)$ to indicate, that the quantities depend on time. The minus sign in F_s and F_d signifies the direction of the forces. This leads to:

$$\begin{aligned} F_s + F_d &= ma(t) \\ -kx(t) - cv(t) &= ma(t) \\ m \frac{d^2x(t)}{dt^2} + c \frac{dx(t)}{dt} + kx(t) &= 0. \end{aligned}$$

This is a **linear differential equation** and it can be solved to find $x(t)$. In the case when the motion is 2D or 3D we need to solve equations of motion for each component of $x(t), y(t)$ and $z(t)$. Consider an electron moving in a uniform (constant) magnetic field perpendicular to its velocity. The force acting on it can be summarized as

$$\begin{cases} \mu \frac{dy(t)}{dt} = \frac{d^2x(t)}{dt^2} \\ \mu \frac{dx(t)}{dt} = \frac{d^2y(t)}{dt^2} \end{cases}$$

for some constant μ .

Deriving Newton's first law

Can we get an understanding of this law by starting from the second law? Let us analyse what happens when $\sum \mathbf{F} = 0$. We acquire

$$\begin{aligned} m\mathbf{a}(t) &= 0 \\ \mathbf{a}(t) &= 0 \quad (\text{for a non-zero mass}) \\ \mathbf{v}(t) &= \mathbf{v}_0 \end{aligned}$$

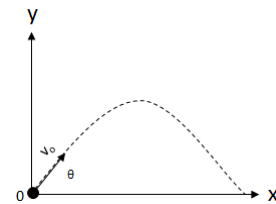
Therefore, we see that the 1st law follows naturally from the 2nd law, i.e. the body will be moving in a **uniform rectilinear motion** or will stay at rest if $\mathbf{v}_0 = 0$.

Exercises

1. A projectile fired at the **origin** at the angle α with an initial speed of v_0 will satisfy the system of equations

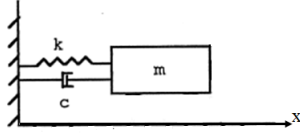
$$\begin{cases} 0 = ma_x(t) \\ -mg = ma_y(t). \end{cases}$$

Solve the system of equations by integrating it twice and applying the right **initial conditions**. By eliminating time from the equation of $y(t)$, show that the trajectory is parabolic.



2. The differential equation governing the motion of a simple pendulum is $\frac{d^2\theta}{dt^2} + \frac{g}{l} \sin(\theta) = 0$, where l is the length of the pendulum, g the gravity constant and θ the angular displacement from the equilibrium position.
 - a) Derive this equation from Newton's second law.
 - b) The derived equation is non-linear and difficult to solve, but, if the angle θ is small ($\theta < 4^\circ$), we can use the following approximation: $\sin(\theta) \approx \theta$. Use this approximation to acquire a linearised equation for the system.
 - c) Find the general solution of the newly derived equation.

3. Find the general solution for a damped oscillator in 1D that is **not** excited by any external forces. What happens when the dampening constant is 0?



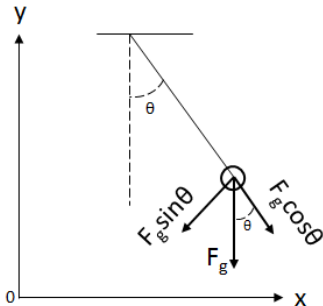
4. A spherical body in free fall in a high-viscosity medium is affected by a drag force $F_d = -cv(t)$. Find the equation of motion in terms of velocity and solve it using the separation of variables.

Answers

1.

$$\begin{aligned} v_x(t) &= v_0 \cos(\alpha) \\ v_y(t) &= -gt + v_0 \sin(\alpha) \\ x(t) &= v_0 \cos(\alpha)t \\ y(t) &= -\frac{gt^2}{2} + v_0 \sin(\alpha)t \\ y(x) &= -\frac{gx^2}{2v_0^2 \cos^2(\alpha)} + \tan(\alpha)x \end{aligned}$$

2. The system:



- a) The force tangential to the bob is $-mg \sin(\theta)$. We have:

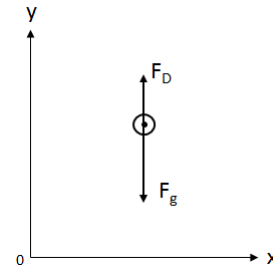
$$\begin{aligned} ma(t) &= -mg \sin(\theta(t)) \\ m \frac{d^2 r(t)}{dt^2} &= -mg \sin(\theta(t)) \\ \frac{d^2 r(t)}{dt^2} &= -g \sin(\theta(t)) \\ \frac{d^2 l \theta(t)}{dt^2} &= -g \sin(\theta(t)) \text{ as } r = l \theta(t) \\ l \frac{d^2 \theta(t)}{dt^2} &= -g \sin(\theta(t)) \text{ as } l \in \mathbb{R}. \\ \frac{d^2 \theta(t)}{dt^2} + \frac{g}{l} \sin(\theta(t)) &= 0 \end{aligned}$$

- b) $\frac{d^2 \theta(t)}{dt^2} + \frac{g}{l} \theta(t) = 0$
 c) The characteristic polynomial is $\lambda^2 + \frac{g}{l} = 0$. That leads to $\lambda = \pm i \sqrt{\frac{g}{l}}$.

What follows is that $x(t) = A \cos(\sqrt{\frac{g}{l}}t) + B \sin(\sqrt{\frac{g}{l}}t)$ or $x(t) = C \sin(\sqrt{\frac{g}{l}}t + \phi_0)$, where $C = \sqrt{A^2 + B^2}$ and $\tan(\phi_0) = \frac{B}{A}$.

3. We have $m \frac{d^2 x(t)}{dt^2} + c \frac{dx(t)}{dt} + kx(t) = 0$, so the characteristic polynomial becomes: $m\lambda^2 + c\lambda + k = 0$. That equation is solved by $\lambda_1 = \frac{-c + \sqrt{c^2 - 4mk}}{2m}$ and $\lambda_2 = \frac{-c - \sqrt{c^2 - 4mk}}{2m}$. So finally $x(t) = A \exp(\lambda_1 t) + B \exp(\lambda_2 t)$. When $c = 0$ we recover the simple harmonic motion with the angular frequency of $\omega = \sqrt{\frac{k}{m}}$.

4. The system:



We have the gravity force $F_g = mg$ and the drag force $F_d = -cv$. That results in

$$\begin{aligned} ma &= F_g + F_d \\ ma &= mg - cv \\ \frac{dv}{dt} &= g - \frac{c}{m}v \\ dt &= \frac{1}{g - \frac{c}{m}v} dv \\ t &= \int \frac{1}{g - \frac{c}{m}v} dv = -\frac{m}{c} \ln \left(g - \frac{c}{m}v \right) + C \\ v(t) &= \frac{gm}{c} - \left(\frac{gm}{c} - v(0) \right) \exp \left(-\frac{c}{m}t \right) \end{aligned}$$